

UNIT-3 WAVE MECHANICS

1. Derive De-Broglie's equation

According to De-Broglie every moving particle is associated with wave. The dual Nature of light possessing both wave and particle properties was explained by combining Planck's expression

for the energy of photon $E = h\nu$ with Einstein's mass energy relation $E = mc^2$ to give

$$mc^2 = h\nu \text{ Or } \nu = \frac{mc^2}{h}$$

$$\frac{c}{\lambda} = \frac{mc^2}{h} \quad (\nu = \frac{c}{\lambda})$$

$$\lambda = \frac{h}{mc} = \frac{h}{p}, \text{ where } \lambda \text{ is the wavelength of photon}$$

$$\text{If particle moves with velocity 'v' then } \lambda = \frac{h}{mv}$$

From above equation if particle velocity is less, then wavelength of wave is more. It was by analogy with this equation associating momentum with a photon that de Broglie expressed the concept of matter wave, according to which a material particle of mass 'm' moving with a velocity 'v' should have an associated wavelength 'λ', called the de-Broglie wavelength.

2. Write difference forms of De-Broglie's Wavelength

(1) De-Broglie's wave equation is given by $\lambda = \frac{h}{mv}$ ---- (1)

If kinetic Energy of particle is $E = \frac{1}{2}mv^2$ or $\sqrt{2Em} = mv$ ---- (2)

substitute Eq (2) in Eq (1)

$$\lambda = \frac{h}{\sqrt{2Em}} \quad (\text{Where E is kinetic Energy})$$

(2) The wavelength associated with a moving particle is independent of any charge associate with it. If the velocity 'v' is given to an electron by accelerating it through a potential difference 'V' then the work done on the electron is Ve. This work done is converted into the kinetic energy of the electron, then

$$eV = \frac{1}{2}mv^2 \text{ Or } \left[\frac{2eV}{m} \right]^{\frac{1}{2}}$$

$$\sqrt{2meV} = mv$$

Substituting this value in the De Broglie equation we have $\lambda = \frac{h}{\sqrt{2meV}}$

Charge of electron is 1.6×10^{-19} C/s and Mass of electron is 9.1×10^{-31} Kg substitute in above equations

$$\lambda = \frac{12.27}{\sqrt{V}} \text{ \AA}$$

3. What are the properties of de Broglie's waves or Matter Waves?

- (1) Lighter is the particle, greater is the wavelength associated with it.
- (2) Smaller is the velocity of the particle greater is the wavelength associated with it.
- (3) When $v = 0$ then $\lambda = \infty$ i.e., wave becomes indeterminate and if $v = \infty$ then $\lambda = 0$ This shows that matter waves are generated by the motion of particles. These waves are produced whether the particles are charged particles or they are uncharged. This fact reveals that these waves are not electromagnetic waves but they are a new kind of waves
- (4) The velocity of matter waves depends on the velocity of material particle i.e. it is not a constant while the velocity of electromagnetic wave is constant.
- (5) The wave nature of matter introduces an uncertainty in the location of the position of the particle because a wave cannot be said exactly at this point or exactly at that point. However, where the wave is large there is good chance of finding the particle while, where the wave is small there is very small chance of finding the particle.

4. Derive Schrödinger time independent equation

Let us consider a group of waves associated with a moving particle. Let ψ represent the displacement of these waves at any time 't'. Let us consider that the wave motion be represent by classical wave equation.

$$\nabla^2 \psi = \frac{1}{v^2} \frac{\partial^2 \psi}{\partial t^2} \text{ ----- (1)}$$

Where v is the velocity. The solution of the above equation is given by

$$\psi = \psi_0 e^{-i\omega t} \text{ ----- (2)}$$

Differentiate above equation with respect to 't'

$$\frac{\partial \psi}{\partial t} = \psi_0 (-i\omega) e^{-i\omega t} \text{ And } \frac{\partial^2 \psi}{\partial t^2} = \psi_0 (-i\omega)^2 e^{-i\omega t} = -\omega^2 \psi$$

$$\frac{1}{v^2}$$

Substitute above value in equation (1) then $\nabla^2 \psi = v^2 (-\omega^2 \psi)$

$$\nabla^2 \psi + \frac{\omega^2}{v^2} \psi = 0 \quad \text{But } \frac{\omega^2}{v^2} = \frac{(2\pi\nu)^2}{(\nu\lambda)^2} = \frac{4\pi^2}{\lambda^2}$$

$$\text{The wave equation is given by } \nabla^2 \psi + \frac{4\pi^2}{\lambda^2} \psi = 0$$

We can substitute the wavelength of the wave accompanying the particle in terms of the particle like property i.e. $\lambda = \frac{h}{mv}$. Then

$$\nabla^2 \psi + \frac{4\pi^2 m^2 v^2}{h^2} \psi = 0$$

If E and V are the total energy and the potential energy of the particle, respectively, then $\frac{1}{2} m v^2 = E - V$ or $mv = \sqrt{2m(E - V)}$

$$\text{The wave equation is given by } \nabla^2 \psi + \frac{8\pi^2 m}{h^2} (E - V) \psi = 0 \quad \text{or} \quad \boxed{\nabla^2 \psi + \frac{2m}{\hbar^2} (E - V) \psi = 0}$$

This equation is known as Schrödinger time independent wave equation.

5. Derive Schrödinger time dependent equation

Let us consider a group of waves associated with a moving particle. Let ψ represent the displacement of these waves at any time 't'. Let us consider that the wave motion be represent by classical wave equation.

$$\nabla^2 \psi = \frac{1}{v^2} \frac{\partial^2 \Psi}{\partial t^2} \quad \text{----- (1)}$$

Where ' v ' is the velocity. The solution of the above equation is given by

$$\psi = \psi_0 e^{-i\omega t} \quad \text{----- (2)}$$

Differentiate above equation with respect to 't' then $\frac{\partial \psi}{\partial t} = \psi_0 (-i\omega) e^{-i\omega t}$

But $\omega = 2\pi\nu$ substitute in above equation $\frac{\partial \psi}{\partial t} = \psi_0 (-i2\pi\nu) e^{-i\omega t} = -2\pi i \frac{E}{h} \psi \quad (E = \nu h)$

$$\text{Or} \quad \frac{\partial \psi}{\partial t} = - \frac{iE}{\hbar} \psi$$

$$\text{Multiplying both side with } \hbar i \frac{\partial}{\partial t} \psi = E \psi$$

Substituting this value in time independent wave equation i.e. $\nabla^2 \psi + \frac{2m}{\hbar^2} (E - V) \psi = 0$

$$\nabla^2 \psi + \frac{2m}{\hbar^2} (i \frac{\partial \psi}{\partial t} - V \psi) = 0$$

$$\frac{2m}{\hbar^2} \nabla^2 \psi + i \frac{\partial \psi}{\partial t} - V \psi = 0 \quad \text{Or} \quad \boxed{i \frac{\partial \psi}{\partial t} = -\frac{2m}{\hbar^2} \nabla^2 \psi + V \psi}$$

Above equation is called as Schrödinger time dependent wave equation. The above equation

can be written as $\boxed{E \psi = H \psi}$ where $i \frac{\partial \psi}{\partial t} = E \psi$ and $H = -\frac{\hbar^2}{2m} \nabla^2 + V$ as a Hamiltonian Operator.

6. Explain Physical significance of wave function (ψ)

The wave function associated with a physical system contains all relevant information about the system and its future behavior and thus describes it completely. It is natural to assume that the wave function be large where the particle is most likely to be and small elsewhere.

If ψ is the amplitude of matter waves at any point in space, then the particle density at that point may be taken as proportional to ψ^2 . Thus ψ is a measure of particle density. When this is multiplied by the charge of the particle, the charge density is obtained. In this way, ψ is a measure of charge density.

According to Max Born $\psi \psi^* = \psi^2$ gives the probability of finding the particle in the state. ψ is a measure of probability density

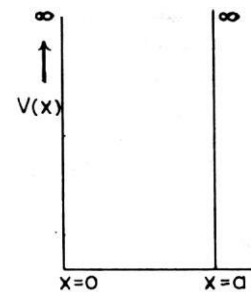
The function $\psi(\mathbf{r}, \mathbf{t})$ is sometimes called probability amplitude of the particle at position \mathbf{r} at time \mathbf{t} . The total probability of finding the particle in the region is of course, unity, i.e. the particle is certainly to be found somewhere in space $\iiint \psi^2 dV = 1$

Limitation of ψ

1. ψ must be finite for all values of x, y, z of the region
2. ψ must be single valued i.e. for each set of values of x, y, z i.e. ψ must be have one value only
3. ψ must be continuous in all region except where potential energy is infinite
4. ψ is analytical i.e. it possesses continuous first order derivative
5. ψ Vanishes at the boundaries.

7. Apply Schrodinger time independent wave equation to particle in Potential Box

Let us consider a square potential well with infinitely high sides, as indicated below fig... If particle in potential well the potential energy is zero. If particle is moving then potential energy increases.



According to Schrödinger time independent wave equation is given by

$$\nabla^2 \psi + \frac{2m}{\hbar^2} (E - V) \psi = 0$$

Potential energy is zero if particle 'x' lies in between 0 to a i.e. $V(x) = 0$ for $0 < x < a$

Boundary condition for wave function is given by

$$\psi = 0 \text{ if particle at } a \text{ i.e. } (\psi)_{x=a} = 0 \text{ And } \psi = 0 \text{ at } 0 \text{ i.e. } (\psi)_{x=0} = 0$$

If potential energy is equal to zero then S. E equation becomes $\nabla^2 \psi + \frac{2m}{\hbar^2} (E) \psi = 0$

$$\text{Or } \nabla^2 \psi + \alpha^2 \psi = 0 \text{ ----- (1) } \left(\alpha^2 = \frac{2mE}{\hbar^2} \right)$$

The solution of above differential equation is $\psi = A \sin \alpha x + B \cos \alpha x$ ----- (1)

Applying the boundary conditions

$$\psi = 0 \text{ at } x = a \text{ and } \psi = 0 \text{ at } x = 0$$

$$\text{i.e. } A \sin \alpha a + B \cos \alpha a = 0 \text{ -----(2)}$$

$$B \cos \alpha a = 0 \text{ -----(3)}$$

If we substitute equation 3 in equation 2 then we have $A \sin \alpha a = 0$

This means wither $A=0$ or $\sin \alpha a = 0$

$$\sin \alpha a = 0 \text{ or } \alpha a = \pi, 2\pi, \dots = n\pi$$

$$\text{Where } n=0, 1, 2, 3, \dots \text{ and } \alpha = \frac{n\pi}{a}$$

If we substitute above condition in equation, we have

$$\psi = A \sin \frac{n\pi x}{a} \text{ ----- (4)}$$

$$\text{And } \alpha^2 = \frac{2mE}{\hbar^2} \text{ or } E = \frac{\alpha^2 \hbar^2}{2m} \text{ or } E_n = \frac{n^2 \hbar^2}{8ma^2}$$

The integer 'n' introduced above is called a quantum number. The E values are called energy levels. The particle that is described by the wave function with a certain n values is said to be in quantum state 'n'. For n=1 the state is called ground state. For 2, 3, ---etc. are known as excited states.

The general form of wavefunction may be written as

$$\psi_n = A \sin \frac{n\pi x}{a} \quad (n = 1, 2, 3, \dots)$$

For A and B values we should normalize above function, normalize condition is

$$\int_0^a \psi^2 dV = 1$$

$$\int_0^a A^2 \sin^2 \frac{n\pi x}{a} dx = 1$$

$$\int_0^a \frac{A^2}{2} (1 - \cos \frac{2n\pi x}{a}) dx = 1$$

$$\frac{A^2}{2} \left[\frac{a}{2\pi n} \sin \frac{2\pi n x}{a} \right]_0^a = 1$$

$$A^2 = \frac{2}{a} \Rightarrow A = \sqrt{\frac{2}{a}}$$

Hence the normalized wavefunctions will have the

$$\text{form } \psi_n = \sqrt{\frac{2}{a}} \sin \frac{n\pi x}{a}$$

The energy level and wave function as shown in the fig.... The Probability of finding the particles more at maximum amplitude.

University Questions

1. Discuss the properties of wave function?
2. What is the de-Broglie's hypothesis? Obtain the De-broglie's wavelength of matter wave
3. Apply Schrödinger's Wave equation to a particle in Infinite Square well potential and obtain wave function and energy values?
4. Write the properties of matter waves and show that matter waves can travel with speed greater than speed of light in vacuum.
5. Obtain Schrödinger's time independent wave equation?
6. Find the K.E and velocity of Photon associated with de-Broglies of wavelength of 0.2865Å ⁰
 $h = 6.6 \times 10^{-34} \text{ Js}, m_p = 1.67 \times 10^{-27} \text{ Kg}$
7. Calculate the wavelength associated with a neutron energy 0.025eV ($m_n=1.67 \times 10^{-27} \text{Kg}$)
8. Compute the de-Broglie wavelength of a proton whose kinetic energy is equal to the rest energy of an electron ($m_e = 9.1 \times 10^{-31} \text{kg}, m_p=1.67 \times 10^{-27} \text{ kg}$)
9. Calculate the energy of an electron wavelength of $3 \times 10^{-2} \text{ m}$. Given $h = 6.62 \times 10^{-34} \text{ Js}$
10. Compute the De-Broglie wavelength of a proton whose kinetic energy is equal to the rest energy of an electron ($m_e = 9.1 \times 10^{-31} \text{Kg}, m_p=1.67 \times 10^{-27} \text{Kg}$) (May06,Dec2004) (2M)
11. Find the first excited state energy of an electron moving along X-axis confined in a box of side length 10^{-10}m (july02) (2M)
12. Calculate the De-Broglie wavelength of an electron which is accelerated by a potential of 100V
 $h = 6.6 \times 10^{-34} \text{ Js}, m = 9.1 \times 10^{-31} \text{ Kg}$
13. Find the energy and momentum of the neutron whose De-Broglie wavelength is 1.5Å ⁰
14. Calculate the de Broglie's wavelength of an electron subjected to a potential difference of 12.5KV .
15. Calculate the energy of an electron wave of wavelength $3 \times 10^{-2} \text{m}$
16. The electron is confined to a box of length 10^{-8} m . Calculate the minimum uncertainty in its velocity.
17. Determine the de-Broglie wavelength of an electron, having kinetic energy of 1eV
18. Compute the energy difference between the first and second quantum state for an electron in one dimension's material having cube side 1m .

1. Write the basic equations of Electricity and Magnetism

$$1. \quad \oint E \cdot dS = \frac{q}{\epsilon_0} \quad \dots(1)$$

This is **Gauss's law of electrostatics** which states that the electric flux through a closed surface is equal to the net charge enclosed by the surface divided by the permittivity constant ϵ_0

$$2. \quad \oint B \cdot dS = 0 \quad \dots(2)$$

This is **Gauss's law of magnetism**. This states that the magnetic flux through a closed surface is zero.

$$3. \quad \oint E \cdot dl = - \frac{d\phi_B}{dt} \quad \dots(3)$$

This is **Faraday's law** of Electromagnetic induction.

This law states that an electric field is produced by changing magnetic field.

$$4. \quad \oint B \cdot dl = \mu_0 i \quad \dots(4)$$

This is **Ampere's law** for magnetic field due to steady current. This law states that the amount of work done in carrying a unit magnetic pole one around a closed arbitrary path linked with the current is μ_0 times the current i .

2. Derive the expression Displacement Current.

According to Basic equation (4) i.e. Ampere's law $\int B \cdot dl = \mu_0 i$ Maxwell's suggested that above equation ' i ' is not total current. He suggested that something must be added in ' i ' of above equations. In order to know this something Maxwell's postulated that similar to the electric field due to changing magnetic field, there would be a magnetic field due to changing electric field. Thus a changing electric field is equivalent to a current which flows as long as the electric field is changing and produces the same magnetic effect as an ordinary conduction current. This is known as displacement current.

And Ampere's law is valid only for steady state phenomena and not for changing fields Let us consider Parallel capacitor for changing field.

Electric field is given by

$$E = \frac{Q}{\epsilon_0 A} \quad (\text{Q is th charge and A is the Area between the plates})$$

For changing electric fields differentiating above equation

$$\frac{\partial E}{\partial t} = \frac{1}{\epsilon_0 A} \frac{\partial Q}{\partial t} = \frac{1}{\epsilon_0 A} i$$

$$i = \epsilon_0 A \frac{\partial E}{\partial t} \quad \left(\frac{\partial Q}{\partial t} = i \right)$$

Displacement Current is given by $i_d = \epsilon_0 A \frac{\partial E}{\partial t}$

Now Modified ampere's law is given by $\oint B \cdot dl = \mu_0 (i + i_d)$

3. Derive Maxwell differential equations

Derivations

$$1. \int E \cdot dS = \frac{q}{\epsilon_0} \quad (\text{Gauss' law for electricity}) \text{ ----- (1)}$$

If ρ be the charge density and dV be the small volume then charge density

$$\rho = \frac{\text{charge}}{\text{Volume}} = \frac{q}{dV}$$

$$q = \rho \times dV$$

For total charge $q = \int_V \rho dV$ Substitute 'q' value in eq.1

$$\therefore \int E \cdot dS = \frac{1}{\epsilon_0} \int_V \rho dV$$

$$\text{or} \quad \int \epsilon_0 E \cdot dS = \int_V \rho dV$$

$$\text{i.e.,} \quad \int D \cdot dS = \int_V \rho dV \quad (\because \epsilon_0 E = D)$$

According to Gauss divergence theorem

$$\oint A \cdot dS = \int_V (\nabla \cdot A) dV$$

$$\text{Hence} \quad \oint D \cdot dS = \int_V (\nabla \cdot D) dV$$

$$\text{so,} \quad \int_V (\nabla \cdot D) dV = \int_V \rho dV$$

$$\therefore \quad \nabla \cdot D = \rho$$

$$\text{or} \quad \nabla \cdot E = \frac{\rho}{\epsilon_0}$$

$$\text{or} \quad \text{div} \cdot E = \frac{\rho}{\epsilon_0} \quad \text{.....(a)}$$

$$\text{or} \quad \frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} = \frac{\rho}{\epsilon_0}$$

$$(2) \quad \oint \mathbf{B} \cdot d\mathbf{S} = 0 \quad (\text{Gauss' law for electricity})$$

Transforming the surface integral into volume integral, we get

$$\oint_S \mathbf{B} \cdot d\mathbf{S} = \int_V \nabla \cdot \mathbf{B} \, dV$$

$$\int_V \nabla \cdot \mathbf{B} \, dV = 0$$

As the volume is arbitrary, the integral must be zero

$$\therefore \quad \nabla \cdot \mathbf{B} = 0 \quad \dots(b)$$

$$\therefore \quad \nabla \cdot \mathbf{B} = 0 \quad \frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} + \frac{\partial B_z}{\partial z} = 0 \quad \dots(b)$$

$$\text{or} \quad \frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} + \frac{\partial B_z}{\partial z} = 0$$

$$(3) \quad \int_S \mathbf{E} \cdot d\mathbf{l} = - \frac{d\phi_B}{dt} = - \frac{\partial}{\partial t} \int_S \mathbf{B} \cdot d\mathbf{S}$$

$$(3) \quad \int_S \mathbf{E} \cdot d\mathbf{l} = - \frac{d\phi_B}{dt} = - \frac{\partial}{\partial t} \int_S \mathbf{B} \cdot d\mathbf{S}$$

$$= - \frac{\partial}{\partial t} \int_S \mathbf{B} \cdot d\mathbf{S}$$

Magnetic flux density is given by

$$B = \frac{\text{Flux}}{\text{Area}(dS)}$$

$$\text{Flux} = B \times \text{Area}(dS)$$

$$\text{Total flux}(\phi_B) = \int_S \mathbf{B} \cdot d\mathbf{S}$$

Substitute ϕ_B in above equation then $\int_S \mathbf{E} \cdot d\mathbf{l} = - \int_S \frac{d\mathbf{B}}{dt} \cdot d\mathbf{S}$

Applying Stoke's theorem

$$\therefore \quad \int_S \mathbf{E} \cdot d\mathbf{l} = \int_S (\nabla \times \mathbf{E}) \cdot d\mathbf{S}$$

$$\frac{\partial \mathbf{B}}{\partial t}$$

As the equation is true is for all surfaces, we have

$$\nabla \times \mathbf{E} = - \frac{\partial \mathbf{B}}{\partial t}$$

$$\text{or} \quad \text{curl } \mathbf{E} = - \frac{\partial \mathbf{B}}{\partial t} \dots(c)$$

$$(4) \quad \oint \mathbf{B} \cdot d\mathbf{l} = \mu_0 I \quad (\text{Ampere's law}) \text{ where } I = i + id$$

$$\text{Let us consider Current density } j = \frac{\text{Current}}{\text{Area}}$$

$$\text{Total Current } I = \oint_S J \cdot dS$$

Substitute above value in Ampere's Law

$$\oint B \cdot dI = \mu_0 \oint_S J \cdot dS$$

Applying Stoke's theorem

$$\oint B \cdot dI = \int_S (\nabla \times B) \cdot dS$$

$$\therefore \int_S (\nabla \times B) \cdot dS = \mu_0 \int_S j \cdot dS$$

$$\text{or } \nabla \times B = \mu_0 j$$

$$\text{But } j = \frac{\text{Total Current}}{A} = \frac{i + i_d}{A} = \frac{i + \epsilon_0 A \frac{\partial E}{\partial t}}{A} = j + \epsilon_0 \frac{\partial E}{\partial t}$$

4. Write a short note on uniform plane wave (Transverse nature)

A uniform plane wave is a particular case of wave equation for which the electric field is independent of y and z and is a function of x and t only. Such a wave is called uniform plane wave.

Consider the case of electromagnetic wave in which the components of vectors E and B vary with one coordinate only (say x) and also with time t, i.e.

$$E = E(x, t) \text{ and } B = B(x, t)$$

$$\text{But } \nabla \cdot E = 0 \text{ or } \frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} = 0$$

$$\therefore \frac{\partial E_x}{\partial x} = 0 \text{ or } E_x = \text{constant} \quad \dots(1)$$

$$\text{Further } \nabla \cdot B = 0 \text{ or } \frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} + \frac{\partial B_z}{\partial z} = 0$$

$$\therefore \frac{\partial B_x}{\partial x} = 0 \text{ or } B_x = \text{constant} \quad \dots(2)$$

Equations. (1) and (2) are obtained on the fact that the derivative of E and B with respect Y and Z are zero.

Further $\text{curl } E = \frac{\partial B}{\partial t}$

$$\therefore \begin{vmatrix} i & j & k \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ E_x & E_y & E_z \end{vmatrix} = - \frac{\partial}{\partial t} [iB_x + jB_y + kB_z]$$

Now $E_x \left[\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} \right] = -i \frac{\partial B_x}{\partial t} = 0 \quad \dots(3)$

$$\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} = 0$$

$$\frac{\partial B_x}{\partial t} = 0 \quad \text{or} \quad B_x = \text{constant} \quad \dots(4)$$

From eq.(3)

Similarly, taking curl B, we can show that E = constant.

Hence, we conclude that E and B are constants as regards to time and space. So these components are static components and hence no part of wave motion. Thus,

$$E = jE_y + kE_z$$

$$B = jB_y + kB_z$$

As vectors E and B do not contain any x-component, and hence I-direction being the direction of propagation of the wave. Further both these vectors are perpendicular to the direction of propagation. **Hence, Maxwell electromagnetic waves are purely transverse in nature.**

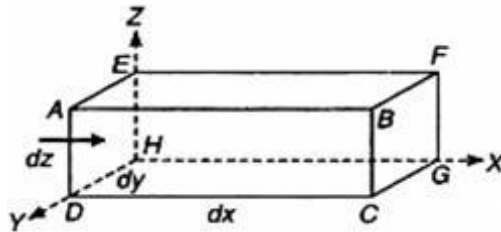
5. State and Derive Poynting vector

One important characteristic of electromagnetic waves is that they transport energy from one point to another point. **The amount of field energy passing through unit area of the surface perpendicular to the direction of propagation of energy is called as Poynting vector.** This is denoted by P. For example in a plane electromagnetic wave; E and B are perpendicular to each other and also to the direction of wave propagation. They involve stored energy. So P has a magnitude $EB \sin 90^\circ = EB$ and points in the direction of wave propagation. The units of P will be $\text{Joule/m}^2 \times \text{sec}$ or Watt/m^2 .

$$P = \frac{1}{\mu_0} (E \times B) \text{ or } (E \times H)$$

Derivation of Expression.

In order to derive the expression for Poynting vector, consider an elementary volume in the form of a rectangular parallelepiped of sides dr , dy and dz as shown in. The volume of parallelepiped is $dx \, dy \, dz$. Suppose the electromagnetic energy



is propagated along the X-axis. Now the area perpendicular to the direction of propagation of energy is $dy dz$. Let the electromagnetic energy in this volume is U . Then the rate of change of energy is

$$\therefore \frac{\partial U}{\partial t} = - \int_S P \cdot dS \quad \dots(1)$$

Negative sign is used to show that energy is entering in the volume. So

$$\int_S P \cdot dS = - \frac{\partial U}{\partial t} \quad \dots(2)$$

We know that:

(1) The energy density per unit volume in electric field E is given by

$$u_E = \frac{1}{2} \epsilon_0 E^2$$

(2) The energy density per unit volume in magnetic field is given by

$$u_B = \frac{1}{2} \mu_0 H^2$$

But $B = \mu_0 H$ and $H = \frac{B}{\mu_0}$

$$u_B = \frac{1}{2\mu_0} B^2$$

Total energy

$$U = u_E + u_B$$

$$\therefore U = \left(\frac{1}{2} \epsilon_0 E^2 + \frac{1}{2\mu_0} B^2 \right) dV$$

The rate of decrease of energy in volume dV is given by

$$- \frac{\partial}{\partial t} \left(\frac{1}{2} \epsilon_0 E^2 + \frac{1}{2\mu_0} B^2 \right) dV$$

Rate of decrease of energy for volume V

$$\begin{aligned} - \frac{\partial U}{\partial t} &= - \frac{\partial}{\partial t} \int_V \left(\frac{1}{2} \epsilon_0 E^2 + \frac{1}{2\mu_0} B^2 \right) dV \\ &= - \int_V \left(\frac{1}{2} \epsilon_0 \frac{\partial E^2}{\partial t} + \frac{1}{2\mu_0} \frac{\partial B^2}{\partial t} \right) dV \\ &= \int_V \left[\epsilon_0 E \left(\frac{\partial E}{\partial t} \right) + \frac{B}{\mu_0} \left(\frac{\partial B}{\partial t} \right) \right] dV \quad \dots(3) \end{aligned}$$

From Maxwell's equations

$$\begin{aligned} \nabla \times \mathbf{E} &= -\frac{\partial \mathbf{B}}{\partial t} \\ (\because \text{Waves are propagating in non-conducting medium}) \\ &= \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \end{aligned}$$

$$\therefore \frac{\partial \mathbf{E}}{\partial t} = \frac{\nabla \times \mathbf{B}}{\mu_0 \epsilon_0} \quad \dots(4)$$

Further $\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$ -----(5)

Substituting the values from eq. (4) and (5) in eq. (3), we get

$$-\frac{\partial \mathbf{U}}{\partial t} = \int_V \left[\epsilon_0 \left(\frac{\nabla \times \mathbf{B}}{\epsilon_0} \right) \cdot \left(\frac{\nabla \times \mathbf{B}}{\epsilon_0 \mu_0} \right) \right] dV \quad \dots(6)$$

Eq.(6) can be written as

$$\begin{aligned} &= \int_V \nabla \cdot (\mathbf{E} \times \mathbf{H}) dV \\ &= \int_V \left[\nabla \cdot (\mathbf{A} \times \mathbf{B}) = \mathbf{B} \cdot (\nabla \times \mathbf{A}) - \mathbf{A} \cdot (\nabla \times \mathbf{B}) \right] dV \end{aligned}$$

Using Gauss theorem of divergence, the volume integral can be expressed in terms of surface integral. Thus,

$$= \int_V (\mathbf{E} \times \mathbf{H}) \cdot \mathbf{n} dS \quad \dots(7)$$

Where n is the unit vector normal to the surface. Comparing eq. (3) with eq. (2), we get

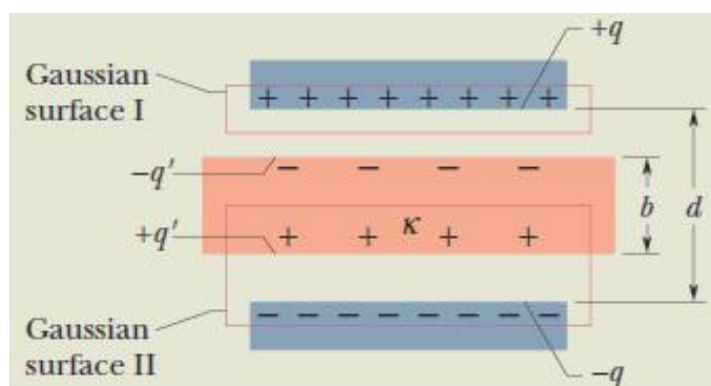
$$\oint_S \mathbf{P} \cdot d\mathbf{S} = \oint_S (\mathbf{E} \times \mathbf{H}) \cdot d\mathbf{S}$$

Or $\mathbf{P} = (\mathbf{E} \times \mathbf{H}) \quad \dots(8)$

In magnitude $P = EH$

This vector shows that energy flow takes place in a direction perpendicular to the plane containing E and H or B. Hence, E and H are the instantaneous values.

6. Derive the relation between D, E and P



Let us consider a parallel plate condenser filled with a dielectric constant 'k'.

When no dielectric present, then Gauss's law is given by

$$\oint E_o \cdot dS = \frac{q}{\epsilon_0}$$

$$E_o \int ds = E_o A = \frac{q}{\epsilon_0} \quad (\text{Where } E_o \text{ is the Electric field without dielectric})$$

$$E_o = \frac{\left(\frac{q}{\epsilon_0 A} \right)}{(1)} \quad ||$$

When dielectric is placed between the plates of the condenser (see fig) the net charge within the Gaussian surface is $q - q^1$. Where q^1 is the induced surface charge. Let E be the resultant field within the dielectric. Then by Gauss's law

$$\int E \cdot dS = \frac{q - q^1}{\epsilon_0}$$

$$E \int ds = EA = \frac{q - q^1}{\epsilon_0}$$

$$E = \left(\frac{q - q^1}{\epsilon_0 A} \right) \quad \text{----- (2)}$$

But dielectric constant $k = E_o / E$ and $E = E_o / K$

Therefore $E = \left(\frac{q}{k \epsilon_0 A} \right)$ (from equation 1)

Substitute E value in equation 2

$$\left(\frac{q}{k \epsilon_0 A} \right) = \left(\frac{q - q^1}{\epsilon_0 A} \right) = \frac{q}{\epsilon_0 A} - \frac{q^1}{\epsilon_0 A}$$

$$\frac{q}{k \epsilon_0 A} = \frac{q}{\epsilon_0 A} - \frac{q^1}{\epsilon_0 A}$$

or $\frac{q}{A} = \epsilon_0 \left(\frac{q}{k \epsilon_0 A} \right) + \frac{q^1}{A}$

We know that $\frac{q}{k \epsilon_0 A} = E$ and $\frac{q^1}{A} = P$

$\therefore \frac{q}{A} = \epsilon_0 E + P$

We put the ratio $\left(\frac{q}{A} \right)$ as D. Hence,

$$\mathbf{D = \epsilon_0 E + P}$$